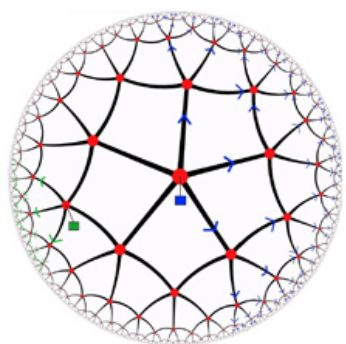


# Quantum Entanglement and the Geometry of Spacetime

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# ENTROPY AND AREA

Bekenstein-Hawking '74:

$$S = \frac{k_B c^3 \text{area}(\text{horizon})}{4G_N \hbar} = k_B \frac{\text{area}(\text{horizon})}{4l_P^2}$$

Planck length  
( $\approx 10^{-33}$  cm)

$G_N \rightarrow$  gravity

$\hbar \rightarrow$  quantum mechanics

$k_B \rightarrow$  statistical mechanics



Event Horizon Telescope '19



What are the “atoms” of the black hole?  
Why is  $S \propto \text{area}$ ?

# ENTROPY AND AREA

If space has  $d$  dimensions,  $G_N \hbar$  has units of area ( $L^{d-1}$ )

*Planck area*: basic unit in quantum gravity, translates into unit of *entropy*

Generalizations of Bekenstein-Hawking:

De Sitter spacetime (Gibbons-Hawking '77):

$$S = \frac{\text{area}(\text{horizon})}{4G_N \hbar}$$

Holographic entropy bounds (Bekenstein '81, Bousso '99):  
for arbitrary closed surface in arbitrary spacetime  $S \leq \frac{\text{area}}{4G_N \hbar}$

Jacobson '95: area-entropy relation *implies* Einstein equation

How general is the area-entropy relation? What is its origin?

A clue: *Holographic entanglement entropy* (Ryu-Takayanagi '06)

Vast (but also limited) generalization of Bekenstein-Hawking

To understand it, we first need to extend our notion of entropy...

# ENTANGLEMENT ENTROPY

Classical mechanics:

definite state  $\rightarrow$  certain outcome for any measurement

Quantum mechanics:

definite state  $\rightarrow$  uncertain outcomes for some measurements

Example:  $|\uparrow\rangle$

measurement of  $S_z$  definitely gives  $+\frac{1}{2}\hbar$

measurement of  $S_x$  gives  $+\frac{1}{2}\hbar$  or  $-\frac{1}{2}\hbar$  with equal probability

When only certain kinds of measurements are allowed, a definite (pure) state will *effectively* be indefinite (mixed)

Suppose a system has two parts, but we can only measure one part

Spin singlet state:  $|0, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle_A |\downarrow\rangle_B - |\downarrow\rangle_A |\uparrow\rangle_B)$   $S_{AB} = 0$

To see that this is a pure state (superposition, not mixture, of  $|\uparrow\rangle_A |\downarrow\rangle_B$  and  $|\downarrow\rangle_A |\uparrow\rangle_B$ ) requires access to both  $A$  and  $B$

For an observer who only sees  $A$ , effective state is mixed:

$$\rho_A = \frac{1}{2} (|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) \quad S_A = \ln 2$$

*Indefiniteness from entanglement*



# ENTANGLEMENT ENTROPY

In general: if  $\rho$  is state of full system, then effective state for subsystem  $A$  is

$$\rho_A = \text{Tr}_{A^c} \rho$$

*Entanglement entropy* is defined as von Neumann entropy of subsystem:

$$S_A = -\text{Tr} \rho_A \ln \rho_A$$

(Any entropy can be viewed as due to entanglement with environment)

Entanglement entropies obey many important properties, such as:

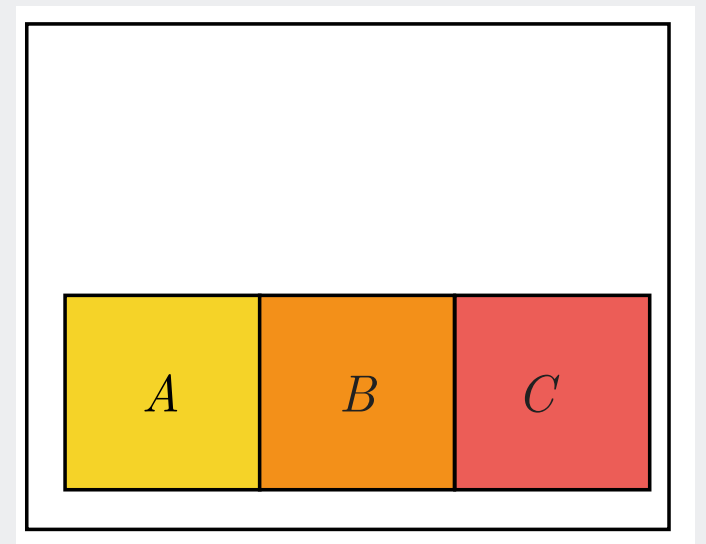
Subadditivity:  $S_{AB} \leq S_A + S_B$

Mutual information:  $I_{A:B} := S_A + S_B - S_{AB} \geq 0$

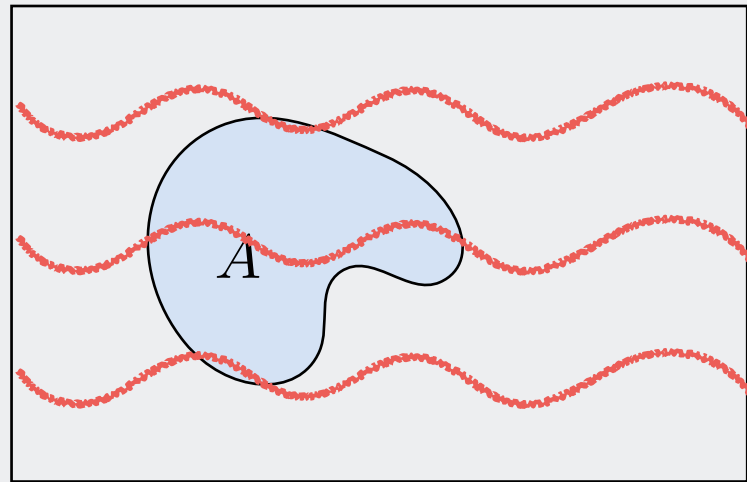
Strong subadditivity:  $S_{AB} + S_{BC} \geq S_B + S_{ABC}$

(Lieb-Ruskai '73)

$$(I_{A:BC} \geq I_{A:B})$$



# ENTANGLEMENT ENTROPY IN QFT



In quantum field theories (& many-body systems), spatial regions are highly entangled with each other

Consider microwave cavity

Even in vacuum, electromagnetic field fluctuates:  
zero-point quantum fluctuations of modes

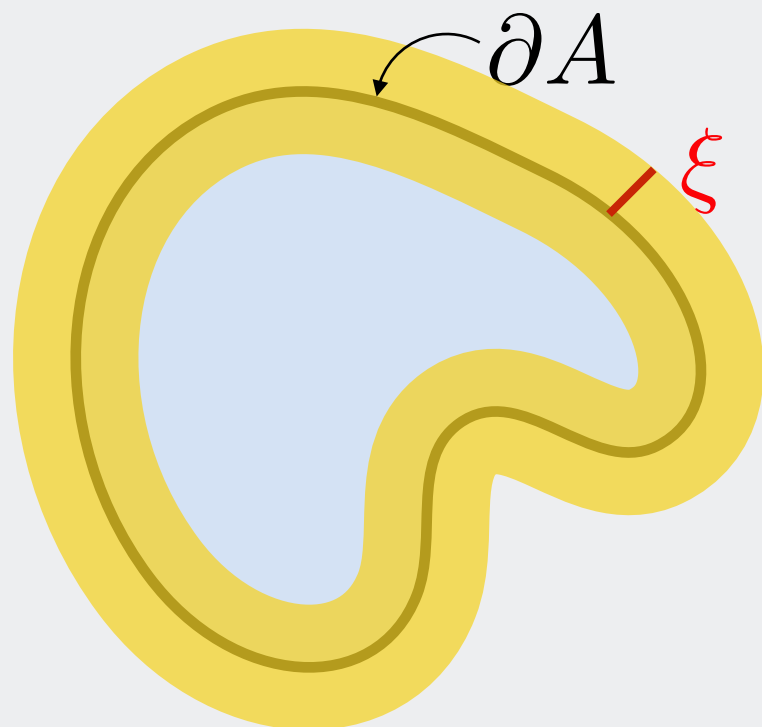
Each mode is distributed in space

=> fluctuations are spatially correlated

=> any part  $A$  of cavity is entangled with rest

$S_A$  is ultraviolet divergent due to entanglement of short-wavelength modes across  $\partial A$

Massive (gapped) field: entanglement extends out to *correlation/Compton length*  $\xi$

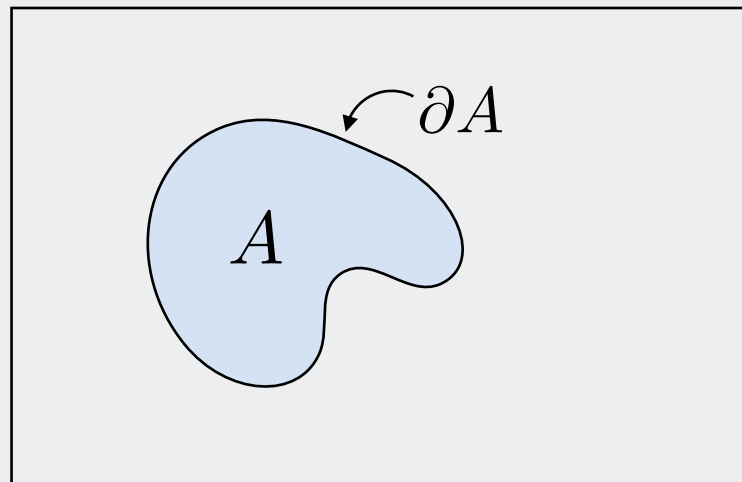


$$S_A = \text{area}(\partial A) \left( \frac{1}{\epsilon^{d-1}} - \frac{1}{\xi^{d-1}} \right) + \dots$$

short-distance cut-off



# ENTANGLEMENT ENTROPY IN QFT



In quantum field theories (& many-body systems), spatial regions are highly entangled with each other

$S_A$  depends on:

- parameters of theory (including  $\epsilon$ )
- state
- size and shape of region  $A$

Contains a lot important physics

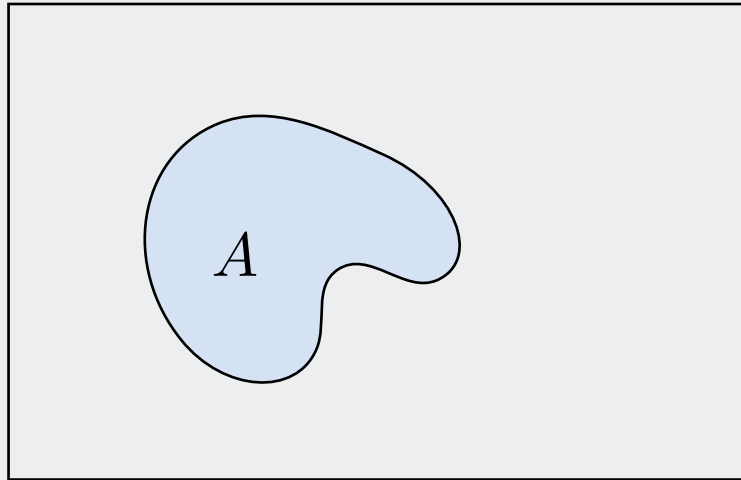
Examples:

- Gapped theory in  $d = 2$ :  $S_A = \frac{L}{\epsilon} - \frac{L}{\xi} - \gamma$ 
  - $\text{length}(\partial A) \curvearrowright \frac{L}{\epsilon}$
  - $\xi$   $\curvearrowright$  correlation length
  - $\gamma$   $\curvearrowright$  topological entanglement entropy (Kitaev-Preskill '05; Levin-Wen '05)

- Critical (conformal) theory in  $d = 1$ :  $S_A = \frac{c}{3} \ln \frac{L}{\epsilon}$ 
  - central charge  $\curvearrowright \frac{c}{3}$
  - $\epsilon$   $\curvearrowright$  short-distance cutoff
  - Diagram: A horizontal line with a blue segment labeled  $A$  of length  $L$ .

- At finite temperature, also usual extensive entropy:  $s(T) \times \text{volume}(A)$ 
  - $s(T)$   $\curvearrowright$  thermal entropy density

# ENTANGLEMENT ENTROPY IN QFT



Powerful probe of QFTs and many-body systems:

- quantum criticality
- topological order
- renormalization-group flows
- energy conditions
- many-body localization
- quenches
- much more...

However, usually very difficult to compute—even in free theories

*Simplifies* in certain theories with *many strongly-interacting* fields...



# HOLOGRAPHIC DUALITIES

Consider a QFT with  $N$  interacting fields

for example  $SU(n)$  Yang-Mills theory,  $N \sim n^2$

When  $N$  is large, these fields may admit a *collective* description in terms of a small number of degrees of freedom

- classical (think of hydrodynamics)
- usually complicated

However, in certain cases, when the fields are very *strongly* interacting, it simplifies dramatically:

General relativity in  $d + 1$  dimensions with cosmological constant  $\Lambda < 0$

(plus some matter fields)

subject to certain boundary conditions: “universe in a box”

(Maldacena '97)

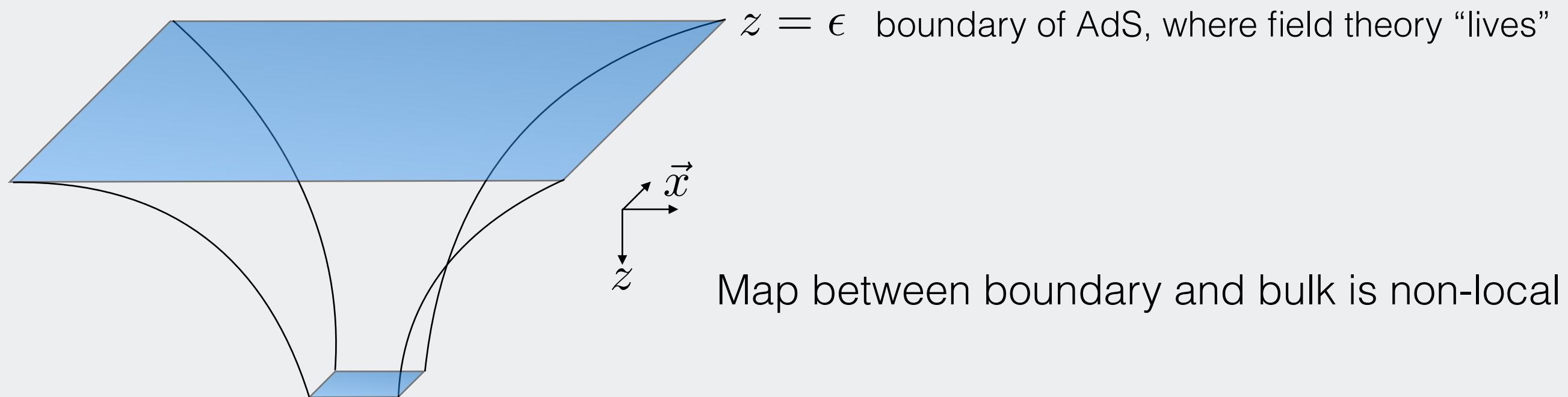
# HOLOGRAPHIC DUALITIES

If QFT is conformal (scale-invariant), ground state is anti-de Sitter (AdS) spacetime:

AdS radius  $d$  dimensions of QFT

$$ds^2 = \frac{R^2}{z^2} (-dt^2 + dz^2 + d\vec{x}^2)$$

extra dimension

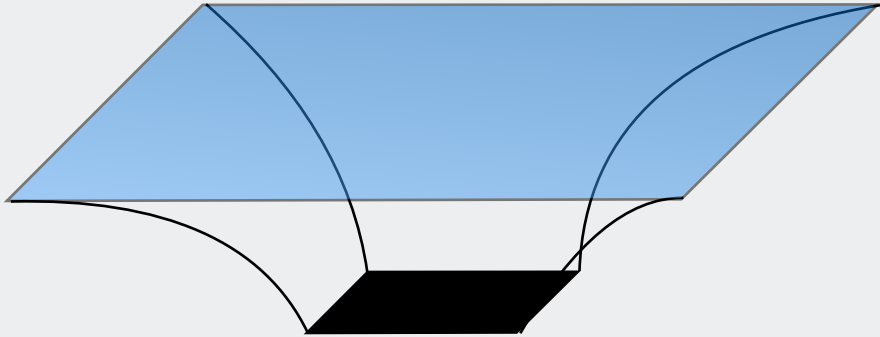


If QFT is gapped (massive), space ends on wall at  $z_{\text{max}} \sim \xi$  (correlation length)

Many specific examples known in various dimensions  
(mostly supersymmetric, derived from string theory)



# HOLOGRAPHIC DUALITIES

QFT	GR
$N$	AdS radius $\curvearrowright R^{d-1}$ $\frac{R^{d-1}}{G_N \hbar} \curvearrowleft$ Planck area
thermodynamic limit $N \rightarrow \infty$	classical limit $\hbar \rightarrow 0$
statistical fluctuations	quantum fluctuations
collective modes	gravitational waves, etc.
deconfined plasma	black hole
$S \propto N$	$S = \frac{\text{area}(\text{horizon})}{4G_N \hbar}$
	
	horizon $\curvearrowright$

Holographic dualities are useful for computing *many* things in strongly interacting QFTs

Let's talk about entanglement entropies...

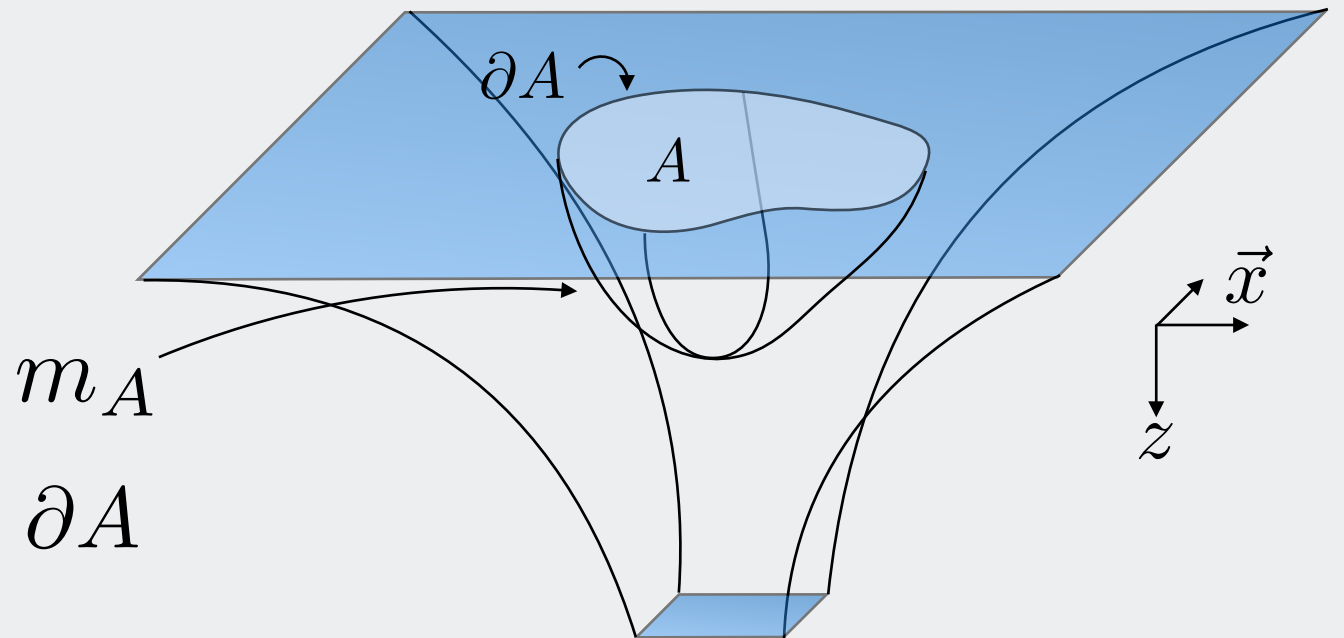
# HOLOGRAPHIC ENTANGLEMENT ENTROPY

Ryu-Takayanagi '06:

$$S_A = \frac{\text{area}(m_A)}{4G_N \hbar}$$

$m_A$  = minimal surface anchored to  $\partial A$

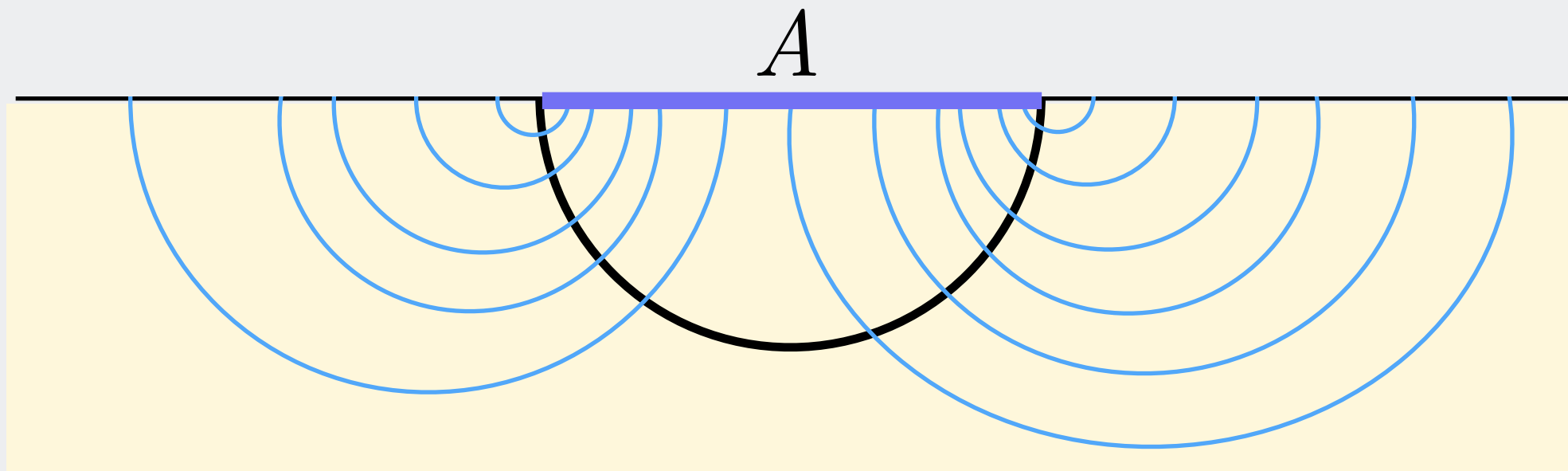
hangs down in order to minimize area



MH-Freedman '16:  $S_A = \max \#$  “bit threads” connecting  $A$  to rest of boundary

Each bit thread has cross section of 4 Planck areas

Represents entangled pair of qubits between  $A$  and complement





# HOLOGRAPHIC ENTANGLEMENT ENTROPY

Ryu-Takayanagi '06:

$$S_A = \frac{\text{area}(m_A)}{4G_N \hbar}$$

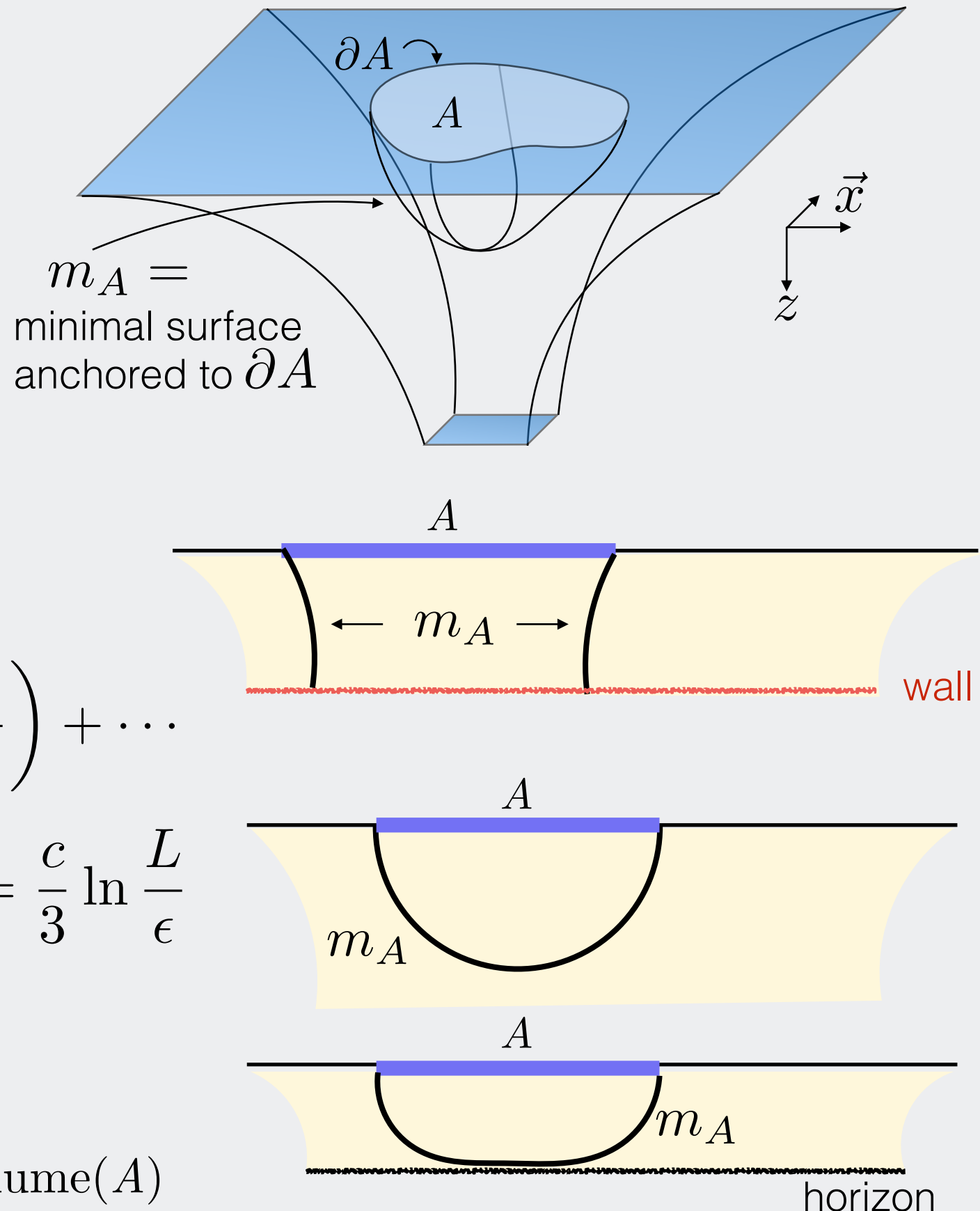
Geometrizes entanglement:

- Area-law UV divergence due to infinite area of  $m_A$  near boundary
- Gapped theory:  
minimal surface extends to wall  
(Klebanov, Kutasov, Murugan '07)

$$S_A = \text{area}(\partial A) \left( \frac{1}{\epsilon^{d-1}} - \frac{1}{\xi^{d-1}} \right) + \dots$$

- Conformal theory in  $d = 1$  :  $S_A = \frac{c}{3} \ln \frac{L}{\epsilon}$

- Finite temperature:  
minimal surface hugs horizon  
=> extensive entropy  $s(T)\text{volume}(A)$



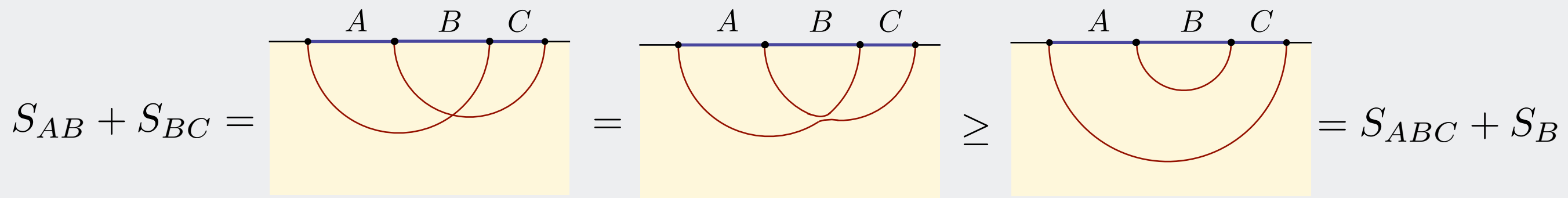
Democratizes Bekenstein-Hawking: not about horizons!

# HOLOGRAPHIC ENTANGLEMENT ENTROPY

*Quantum information theory is built into classical spacetime geometry*

Example: Strong subadditivity

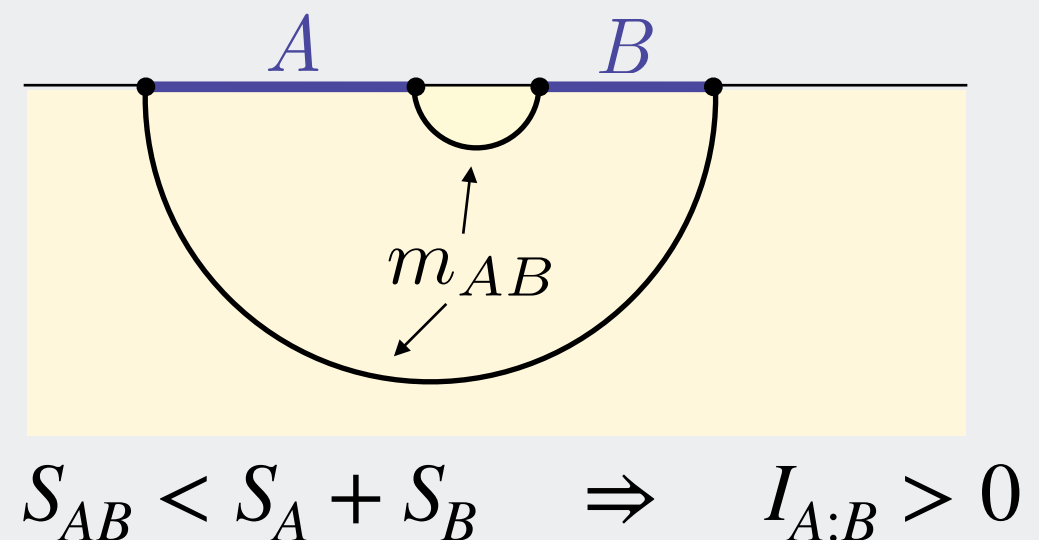
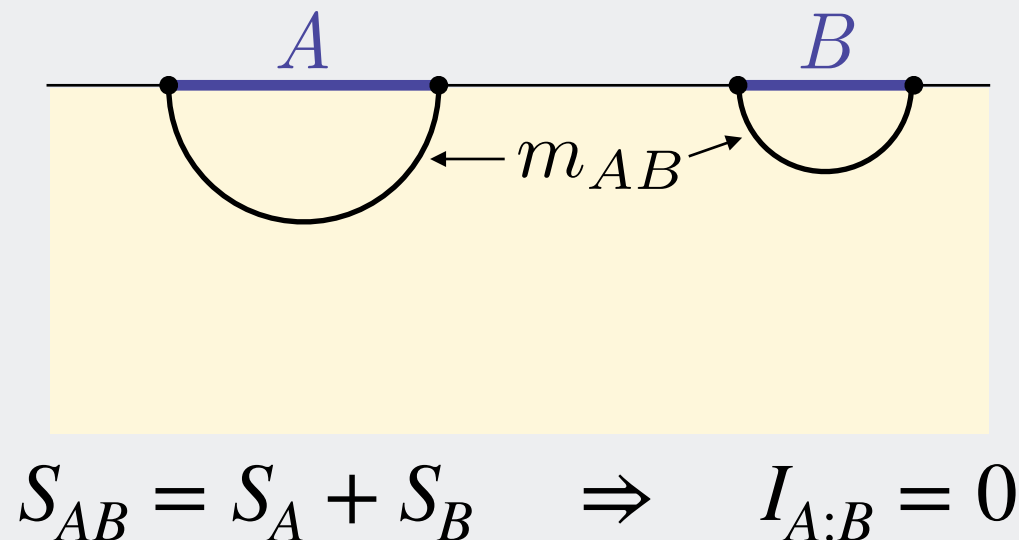
$$S_{AB} + S_{BC} \geq S_{ABC} + S_B$$



(MH-Takayanagi '07)

In fact, RT formula obeys *all* general properties of entanglement entropies  
(Hayden-MH-Maloney '11; MH '13)

Also has special properties, such as phase transitions (MH '10)



# HOLOGRAPHIC ENTANGLEMENT ENTROPY

Many other developments:

- Many discoveries about entanglement in field theories using holography
- Time dependence ([Hubeny-Rangamani-Takayanagi '07](#))
- Derivation of RT formula ([MH '10](#), [Lewkowycz-Maldacena '13](#))
- Einstein equation from RT ([Lashkari-McDermott-Van Raamsdonk '13](#))
- Quantum ( $1/N$ ) corrections ([Faulkner-Lewkowycz-Maldacena '13](#))
- Special properties obeyed by holographic entropies ([Bao et al '15](#))
- Tensor networks for modelling holography ([Swingle '08](#))
- Bit threads & entanglement structures ([Cui et al '18](#))
- Subregion duality ([MH-Hubeny-Lawrence-Rangamani '14](#))
- Holography as quantum error-correcting code ([Almheiri-Dong-Harlow '14](#))
- Other quantities: relative entropies, entanglement of purification, ...
- ...

The big one:

- Does space emerge from entanglement? From bit threads?